

## Notes 10 : Intermittency, Fractional Brownian Motion, Hurst Parameter

→ Here : Consider  $\rightarrow$  FBM (Fractional / Fractal Brownian Motion)

$\left\{ \begin{array}{l} \rightarrow \text{Hurst Parameter} \\ \rightarrow \text{Temporal Intermittency.} \end{array} \right.$

**Fractal character T.B.D.**

### - Harold E. Hurst's Story

- hydrological engineer
- meticulous record keeper, observer
- active in Aswan Dam's construction

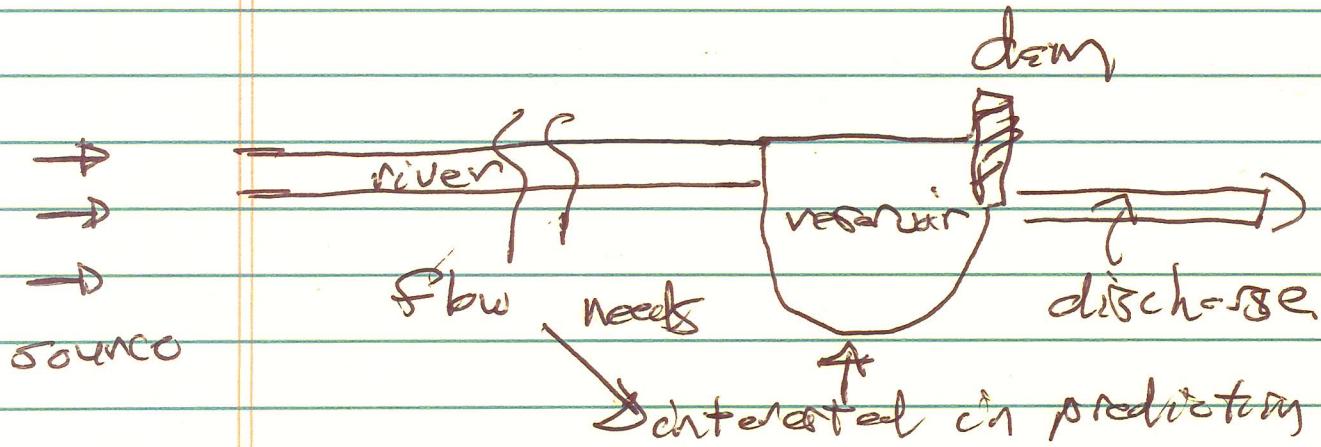
### Problem:

→ statical description of flow, discharge  
of Nile

→ prediction, modelling based on time  
series

i.e. how big should the Aswan  
reservoir be?

i.e schematic:



Challenge:

- characterize the time variation of the river discharge.
  - how big a reservoir?
- analogous to machine construction.

Data - Time Series.

- For long time average in stationary state, might expect random walk, etc

i.e.  $B_H(t)$  → same general, stationary  
 (i.e. river flow due to precipitation) time series.  
 Then expect:

$$E \left\{ (B_H(t+T) - B_H(t))^2 \right\} = T^{(1)}$$

but instead get:

- periods of sustained precipitation, discharge:  
 → "Joseph effects" - persistence

c.f. as in ~~"7 years of feast, famine"~~, in Bible

- "Noah effect" (large outliers)  
 "Black Swan")
- anomalously large event / flood

c.f. Great Flood

- see:
- Mandelbrot and Wallis 1968  
 + technology.
  - Mandelbrot and Van Ness

i.e. not random, instead BM.

$$E \left\{ (B_H(t+\tau) - B_H(t))^2 \right\} = \tau^{2H}$$

$H \rightarrow$  Hurst exponent  
Holder

$H = 1/2 \rightarrow$  Brownian Motion

instead find  $0 < H < 1$

$\rightarrow 1/2 < H < 1$

- memory, positive correlation

- long term persistence

- "Joseph" - avalanche

- super-diffusion

$$\text{i.e. } H = 1/2$$

$$\langle \delta B^2 \rangle \sim \tau^2$$

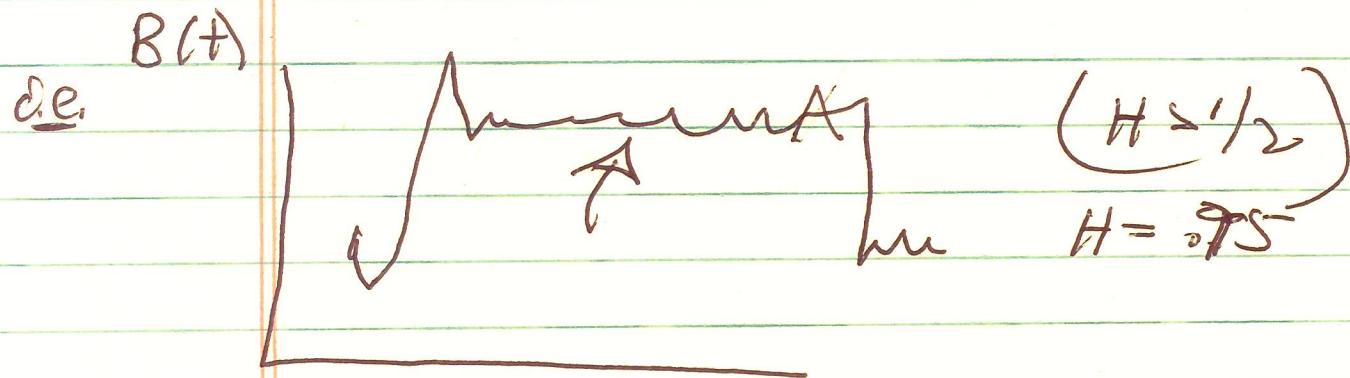
$\Leftrightarrow$

"ballistic" "pulses"

$$\rightarrow 0 < H < \frac{1}{2}$$

- temporal anti-correlation (weak excursion)  $\rightarrow$  does not much excursion
  - $H_1/\text{low}$  value switching
  - sticking
  - sub-diffusion
- $\Delta B \sim T^\alpha$   
 $\alpha < \frac{1}{2}$

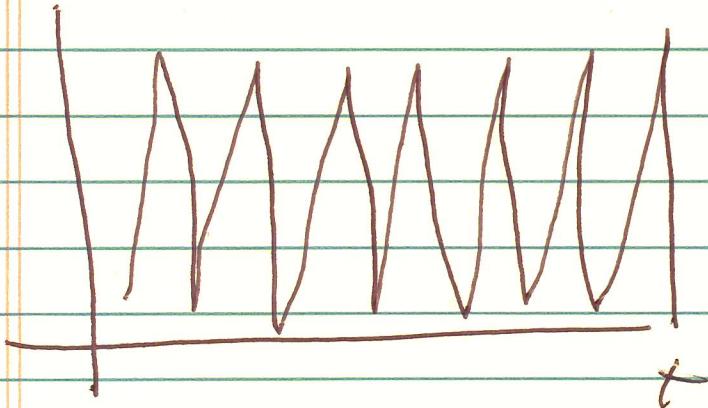
$\rightarrow H = \frac{1}{2}$ , Wiener Process  
Brownian Motion



long term persistence ..  
pulse  $\rightarrow$  auto-similarity

$\beta(t)$  $(H < \frac{1}{2})$ 

$$H \sim 0.04$$



Cycling  $\rightarrow$  anti-persistent

Point:  $H$  measured represents memory in dynamics.

(N.B., Recall:  $\Omega$  measures  $f/f_0$  scaling (in dep  $E$ )  $\leftrightarrow$  memory.)

Note: Can re-write as:

$$E \{ (AB)^2 \} = (\Delta t)^{2H}$$

so

$$\boxed{\ln |AB| / \ln |\Delta t| = H}$$

Second recall 2

$$\frac{D_B}{b} = \ln N / \ln (1/\epsilon)$$

box counting dimension

similarity

$H \leftrightarrow D_B$  (hence Holder exponent).

So  $H$  is (obviously) related to dimensionality of "Brownian" process

→ hence "Fractional Brownian Motion"  
(generalization).

- Generalization of B.M. can't account for "Noah, Joseph phenomena".

A bit more:

→ Hurst observed empirically the ideal reservoir capacity  
(i.e. big enough to hold "biggest!")

8

seek characterize, predict

d.e.  $R(\delta) \equiv$  ideal capacity

$S(\delta) \equiv$  standard deviation of discharge

$\delta \equiv$  # of successive discharge

$$R(\delta)/S(\delta) \sim \delta^H$$

d.e.  $\begin{cases} R \rightarrow \text{pile content} \\ S \rightarrow \text{standard Deviation} \\ \delta \rightarrow \text{time} \end{cases}$

empirically,  $H \sim .7 \rightarrow .85$

for Nole  
(Austria)

(significant deviation from B.M.)

River Flow  $\rightarrow$  series of pulses.  
(see 8e.)

$\rightarrow$  How to  $H$ ?

Consider a time series:

$x_1, x_2, \dots, x_n$

N.B. Can apply Kinematic Wave theory to problem of dynamics of river

i.e. traffic:  $\frac{\partial \rho}{\partial t} + \frac{\partial \xi}{\partial x} = 0$

cf whitham

$$\xi = Q(\rho)$$

$$\frac{\partial_t \rho}{\partial t} + c(\rho) \frac{\partial \rho}{\partial x} = 0$$

$$Q(\rho)$$



etc.

Then, for floods:

$$\Delta \rightarrow A(x,t) \rightarrow \text{cross sectional area of river-bed.}$$

$$\frac{\partial A}{\partial t} + \frac{\partial Q}{\partial x} = 0$$

$$Q = Q(A, x)$$

$$\underline{\underline{0}} \quad \frac{\partial}{\partial t} A + \frac{\partial Q}{\partial A} \frac{\partial A}{\partial x} = - \frac{\partial Q}{\partial x}$$

$$C = \frac{\partial Q}{\partial A} = \frac{1}{b} \frac{\partial Q}{\partial h}$$

width      ht.

and can use relations (empirical)

$$V = \Sigma / A$$

$$\left\{ \begin{array}{l} V = \left( \frac{A \sin \alpha}{P} \right)^{1/2} \\ Q = V A \end{array} \right.$$

wetted perimeter  
↓

i.e.  $F_g = C_f P V^2 \rho_0$

$\downarrow$   
coeff

$$\text{and } F_g = \rho g A \sin \alpha$$

Balance  $\rightarrow V$ .

$$\text{so } Q - A^{3/2} \text{ etc.}$$

$$\Rightarrow \partial_t A + C(A) A_x = 0$$

etc

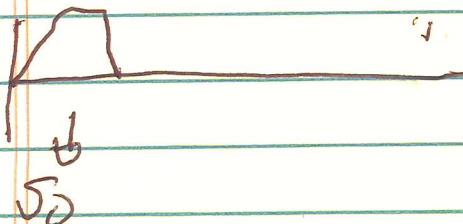
Can extend to:

$$\left\{ \begin{array}{l} \partial_t h + v \partial_x h + h v_x = 0 \\ \partial_t v + v \partial_x v + g' h_x = g' S - \frac{Cv^2}{h} \end{array} \right.$$

etc.  $\rightarrow$  similar to dynamic traffic flow,  
and shallow water.

then, consider localized, noisy source  
in  $h$  equation

$$\partial_t h + v \partial_x h + h v_x = S_0(x, t)$$



$$S_0 = \mathcal{T}d(t-t') F_0(x)$$

etc.

8to

What is/are  $\{ \text{correlation, } H \\ \text{statistics, } \dots \}$   
of downstream flow, ht, etc. ??

$$P_t : NL + N_{\text{loss}} \sim$$

Flooding ?

What is sensitivity to excitation ?

then:

$$Cn^H = \underset{\text{expectation}}{E} \left[ \frac{R(n)}{S(n)} \right]$$

$R(n)$  = Range of first  $n$  values

$S(n)$  = std Deviation first  $n$ .

More quantitatively:

1) mean,  $m = \frac{1}{n} \sum_{i=1}^n x_i$

2) adjust series to mean, (de-mean).

$$y_t = x_t - m \quad j = 1, \dots, n$$

3) calculate cumulative deviate from mean

$$z_t = \sum_{i=1}^t y_i$$

4) Compute Range of deviate:

$$R(n) = \max(Z_1, \dots, Z_n) - \min(Z_1, \dots, Z_n)$$

5) Computed std. Dev.

$$\sigma(n) = \left( \frac{1}{n} \sum_{i=1}^n (x_i - \bar{x})^2 \right)^{1/2}$$

~~$R(n)/\sigma(n)$~~  goes  $\rightarrow H$  says over  
partial series

i.e.  $H = \ln \left[ \overline{\left[ R(n)/\sigma(n) \right]} \right] / \ln(n)$

-  $H \rightarrow$  "R/S analysis"

- Key:  $R(n) \rightarrow$  Range

Masures 'dispersion' in time series.

Some similarity to Gini Parameter  
in Econ.

- Gini Coeff - Measures dispersion / Concentration of wealth in Population

For wealth  $x_{c^*j}$

$$G = \frac{1}{n} \sum_{i=1}^n \sum_{j=1}^n |x_i - x_j| / 2n \sum_{i=1}^n x_i$$

$\uparrow$   
dispersion

$G=1$  if 1 person has all  $\$$ .

Some similarity R and G.

→ Gini or analogue can be useful as measure of concentration in a time series.

Details re H:

- can ask, is high H (Euler) randomness a mesurement for  $D(t)$  → fine very diffusing.
- Remedy  $N$  series  
divided into shorter  
 $n = N, N/2, N/4, \dots$

then rescaled range calculated for each  $n$ .

### Related Issues

- $B(f)$  self-similar / self-affine  
with Hölder (Hurst) dimension  $H$   
if
- Can parallel fractal / multi-fractal:

$$\left\{ E[(B(t+\tau) - B(t))^2] \right\}^{1/2} \sim \text{const } \tau^H$$

→ unfractal scaling  
(applies all  $\tau$  moments)

VS.

$$\left\{ E[(B(t+\tau) - B(t))^2] \right\}^{1/2} \sim \tau^{H(\epsilon)}$$

c.r. dependent on  $\epsilon$

⇒ multi-fractal

→  $H$  defines "roughness" of series  
 " + "  
 "randomness"

→ Can define  $\langle B^2(\omega) \rangle \approx B_\omega^2 = \int_t^{t+T} \langle B(t) B(t+\tau) \rangle$   
 Result:

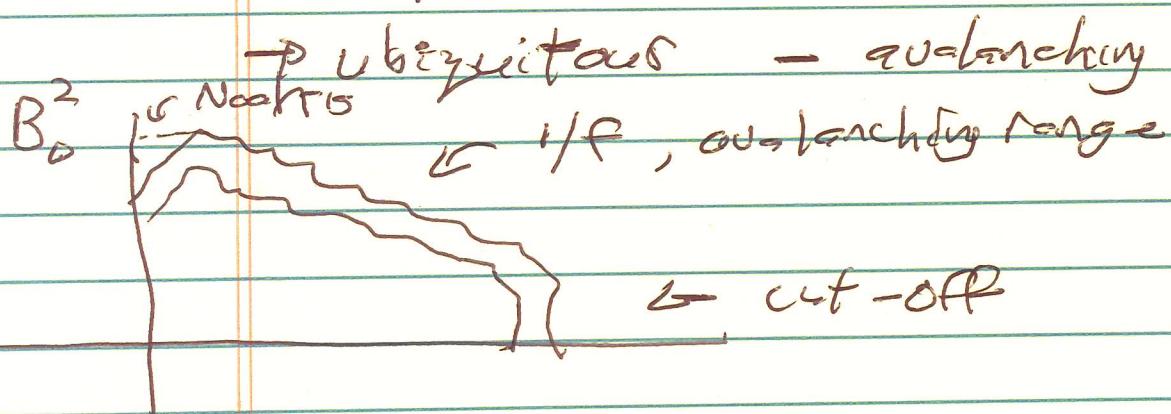
$$\langle B^2 \rangle_\omega \approx \omega^{-B}$$

$$B = 2H - 1$$

$$\sim H=1 \quad \langle B^2 \rangle_\omega \sim 1/\omega$$

" 1/f noise "

→ Power at low frequencies - persistence



$$\sim H=1/2 \quad (BM)$$

$$\langle B^2 \rangle_\omega \sim 1 \quad \text{white noise}$$

$$\sim H = 0$$

$$\langle B^2 \rangle_\omega \sim \omega \quad \rightarrow \text{coherent}$$

$\rightarrow$  Can define Fractal dimension of time series:

$$D = 2 - H$$

i.e.  $H=0$   $D=2$

$$H=1 \quad D=1$$

$\rightarrow H \sim \frac{1}{2} \rightarrow$  "mild"

$H \sim 1 \rightarrow$  "wild" large variation.

# Mild Variation

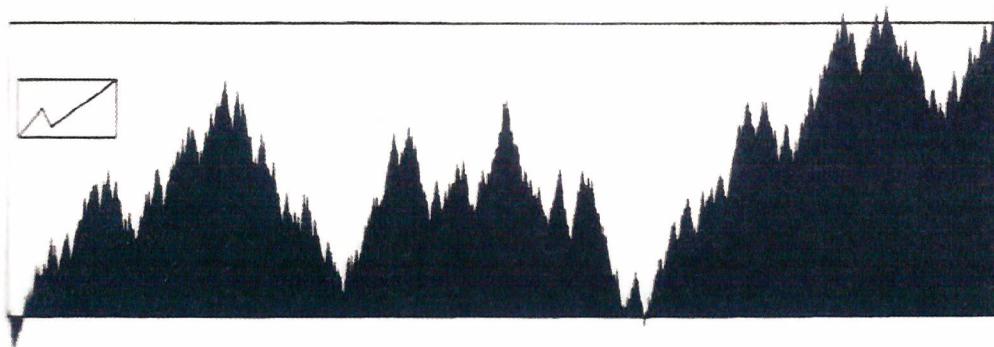
15.

## Wiener Process



E6 ◊◊ ...PANORAMA OF SELF-AFFINE VARIABILITY

183



Wiener  
cumulative



increment

~~generator~~

→ White  
noise



FIGURE E6-6. The top line illustrates a cartoon of Wiener Brownian motion carried to many recursion steps. The generator, shown in a small window, is identical to the generator A2 of Figure 2. At each step, the three intervals of the generator are shuffled at random; it follows that, after a few stages, no trace of a grid remains visible to the naked eye.

→ The second line shows the corresponding increments over successive small intervals of time. This is for all practical purposes a diagram of Gaussian "white noise" as shown in Figure 3 of Chapter E1.

Mandelbrot  
on  
Finance

# Wild Variation

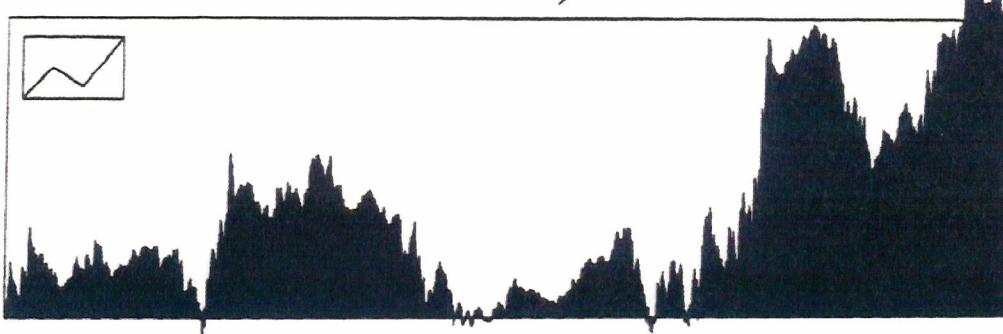
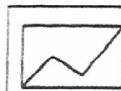
16.

184

SELF-SIMILARITY AND... ◊ ◊ E6

Wild Variation

Wiener process  
multi-fractal  
trading time



increments  
non-Gaussian  
serial dependence  
high variability  
of increments  
→ wild

diffn good

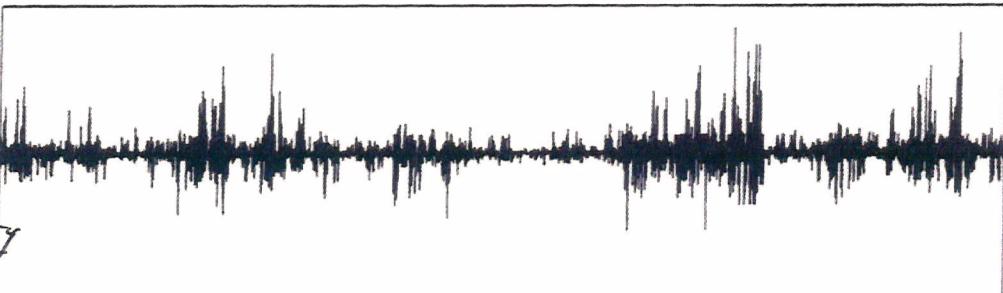


FIGURE E6-7. This figure reveals – at long last – the construction of Figure 2 of Chapter E1. The top line illustrates a cartoon of Wiener Brownian motion followed in a multifractal trading time. Starting with the three-box generator used in Figure 6, the box heights are preserved, so that  $D_T$  is left unchanged at  $D_T = 2$  (a signature of Brownian motion), but the box widths are modified. (Unfortunately, the seed is not the same as in Figure 6.)

The middle line shows the corresponding increments. Very surprisingly, this sequence is a “white noise,” but it is extremely far from being Gaussian. In fact, serial dependence is conspicuously high. The bottom line repeats the middle one, but with a different “pseudo-random” seed. The goal is to demonstrate once again the very high level of sample variability that is characteristic of wildly varying functions.

The resemblance to actual records exemplified by Figure 1 of Chapter E1 can be improved by “fine-tuning” the generator.

$\rightarrow 1/f$  noise ( $\langle 1/f^x \rangle \propto \omega^{-1}$ )

$\sim 1/f$  noise is ubiquitous in physical systems

$\sim$  was a driver for SOC theory

$\sim$  low frequency power suggestive of larger clusters (percolation clusters)

$\sim$  suggestion of scale invariance

$\rightarrow$  Getting  $1/f$ .

- not easy, i.e. unless:

$$\langle \hat{\phi}(t_1) \hat{\phi}(t_2) \rangle = |\phi|^2 e^{-(|t_2 - t_1|/\tau_0)}$$

$$\Rightarrow S(\omega) = \frac{1/\tau_0}{\omega^2 + 1/\tau_0^2} \sim 1/\omega^2$$

i.e.  $1/\tau_0$  imposes scale, but  $1/f$  noise scale free (self-affinity).

- if conserved order parameter

$$\partial_t \phi + \nabla \phi \cdot \nabla \phi = 0$$

$$\langle \phi(t) \phi(\tau) \rangle_k = |\phi|^2 e^{-k^2 D \tau}$$

(Follows Taylor +  
McNamee)

$$S(\omega, k) = |\phi|^2 \frac{k^2 D}{\omega^2 + (k^2 D)^2}$$

$$T_c \rightarrow \infty \quad k \rightarrow 0.$$

Can recover order invariance.

- Alternative: Montroll (postel)

- now consider ensemble of random processes, each with  $T_c$   
(i.e. distribution), probability  $P_c$ .

$$S(\omega)_{\text{eff}} = \sum_{T_c} P(T_c) S_{T_c}(\omega) dT_c$$

- Demand  $P(T_c)$  scale invariant

$$P(T_c) = 1/T_c \quad (\text{dim})$$

$$\Im(\omega) = \tan^{-1} \frac{\omega T_c}{\omega} \Big|_{T_c}$$

$$\sim 1/\omega \Rightarrow 1/f$$

N.B.:  $P \approx 1/T_c$

- long events rare

- short events numerous

$\Rightarrow$  Zipf's Law

$$P(\Delta x) \sim 1/(\Delta x) \Leftrightarrow 1/f$$

N.B. Lognormal well approximated by  
 $P \propto 1/x$  over finite range

Guided R/S  $\leftrightarrow$  Hurst exponents A Broeder P.

Now  $\rightarrow$  Characterizing "wild" - Vision of Ats

$\rightarrow$  How characterize "wild" randomness?

$\rightarrow$  non. distribution  
- Levy flights are prime example of wild randomness

- Levy Flight (pioneered by Paul Levy)  
Levy process (Corset - Mandelbrot)

is random walk in which  $\Delta X$  distributed along  $P(\Delta X)$  where  $P(\Delta X)$  has "heavy" tail ( $\rightarrow$  power law).

e.g. Cauchy flight,  $P(u) \sim A/u^2$

- Specific example: probability distribution probability

Consider  $P(T) = P(T > u)$   
step size

$$P(T > u) = \begin{cases} t^{-\frac{1}{\alpha}} & : u < t \\ u^{-\frac{1}{\alpha}} & : u \geq t \end{cases}$$

power law

derived from Pareto distribution of incomes (power law).

More generally:  
density

$$\{ P(C > u) = O(u^{-k}) \}$$

$$1 < k < 3$$

Brought to:

OV:  $\{$  Pareto-Levy Law  
Mandelbrot 1960.

→ emerged from economics, concerned with income distributions, especially tail.

→ Pareto (1897)  
Levy (1925)  $\} \quad \begin{aligned} &\rightarrow \text{observed power law} \\ &(P \propto L) \Rightarrow \begin{cases} \text{Wild flights} \\ \text{Levy} \end{cases} \\ &\rightarrow \text{noted that P-L} \\ &\text{distribution satisfies} \\ &\text{a Limit Theorem} \\ &(\text{but not Gaussian}) \end{aligned}$

→ Strong Pareto Law:

$P(u) \equiv \% \text{ of indiv. with income } U > u_0$

$$P(u) = \begin{cases} (u/u_0)^{-k}, & u > u_0 \\ 1, & u < u_0 \end{cases}$$

## Power Laws

then density  $\rightarrow p(u) = -dP(u)/du$  :

(pdf)

$\rightarrow$  Power Law.

$$p(u) = \begin{cases} \alpha u^{-\alpha-1} & u > u^* \\ 0 & u \leq u^* \end{cases}$$

$p(u)$  characterized by  $\left. \begin{array}{l} u^* \rightarrow \text{scale factor} \\ \alpha \rightarrow \text{inequality index} \end{array} \right\}$

$P(u)$  fits broad range of populations  
 (US tax payers, residence towns etc.)  
 (debated)  $\left( \text{chart of robustness} \right) \Rightarrow$   $\left. \begin{array}{l} \text{pdf is} \\ \text{attractor in} \\ \text{fitu space} \end{array} \right\}$

$\rightarrow$  Weak Pareto Law  $\ddagger$  (more robust)

$$\begin{aligned} P(u) &\stackrel{\text{"behaves like"}}{\sim} (u/u_0)^{-\alpha} \quad u \rightarrow \infty \\ \Rightarrow p(u) &\approx (u/u_0)^{-\alpha-1}, \quad \left. \begin{array}{l} \alpha < 2 \Rightarrow \\ \text{fat tail} \end{array} \right\} \\ \left. \begin{array}{l} \text{need } \alpha > 2 \\ \text{for 2nd moment} \\ \text{convergence} \end{array} \right\} \end{aligned}$$

N.B. Competitors for Pareto:

- exponential tail:  $\rightarrow ?$

$$p(u) = \kappa u^{-\alpha-1} e^{-bu} \quad \left. \begin{array}{l} b \rightarrow 0 \\ \approx \end{array} \right\}$$

20

- log-normal (why log normal relevant to income)

## $\Rightarrow$ Thermodynamic Theories (?)

- noting that Gaussian arises from Brownian motion  $\Rightarrow$  many small kicks in velocity,

ask

- can economic interactions exchange characteristics of money leading to  $P-L$  in equilibrium?  
 $P_S P-L$  result of ~~?~~ a limit  
 Theorem B

$\Rightarrow$  NO! / Yes!

large  
events  
unstable

- $P(u)$  decreases too slowly, large  $U$ .

- might try  $\ln U = V \Rightarrow$  heads to lognormal (can speak of additivity of  $\ln U$  increments, and convergence),

all debatable

Percolation  $\rightarrow$  small  $\rightarrow$  wild ?!

but

23

## Pareto - Levy Random Variables

- Issue: Pareto law resilient to how income computed!
- $\Rightarrow$
- Law emerges as a Limit Theorem.

i.e.

### Levy Stable Distributions

(attractors  
in  
phase space)

- $U_i \rightarrow$  statistically indep. incomes (ref to scale, origin)
- $U', U''$  follow P-L, then:  
 $U' + U''$  follows law

i.e. addition of random variable, linear comb. on PL

No random coeff

$$(a' U' + b') + (a'' U'' + b'') = a U + b$$
$$a', a'' > 0, b, b'' > 0 \Rightarrow a > 0, b.$$

two.

i.e. adding to P-L Law incomes  $\Rightarrow$  income "on" P-L Law.

$\therefore \{$  P-L law is an example of an L-stable process  $\}$

$\rightarrow \left\{ \begin{array}{l} P_L \text{ densities } 23. \\ \alpha = 1.2, 1.5, 1.8 \end{array} \right.$

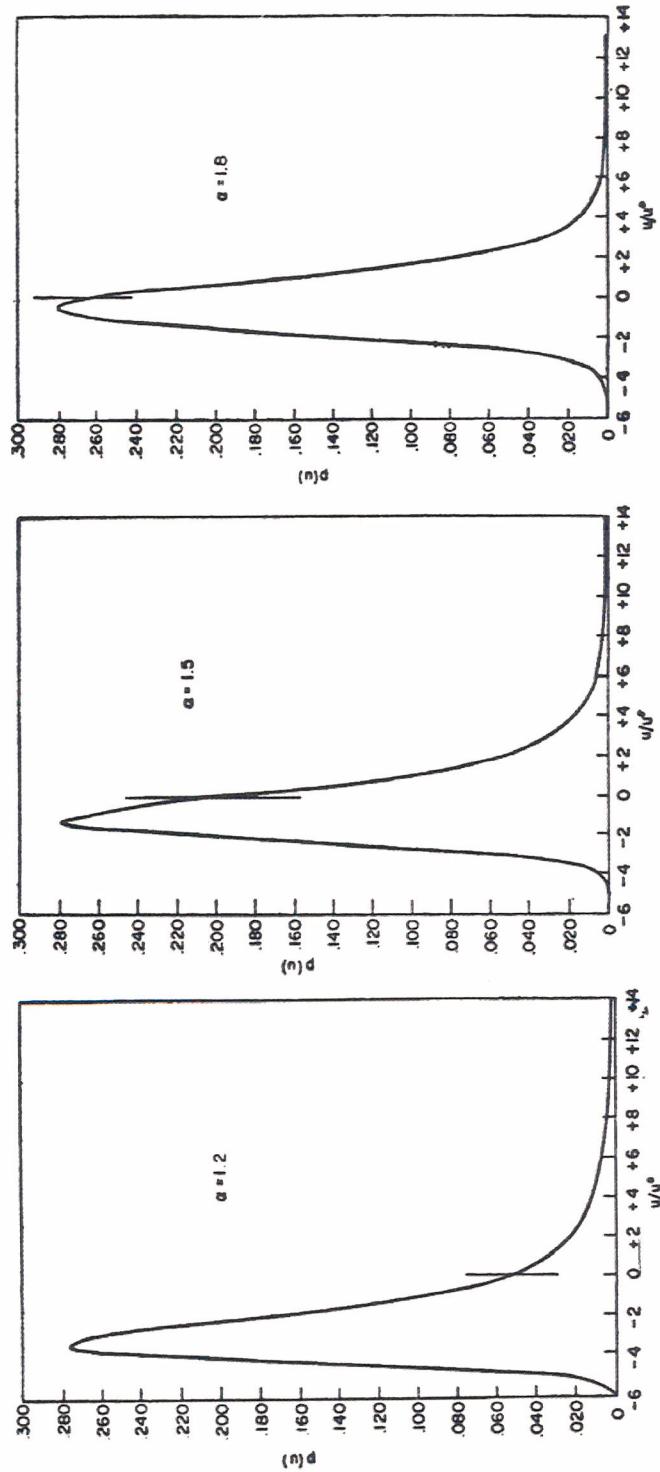


FIGURE 1: DENSITIES OF REDUCED P-L VARIABLES, FOR  $M=0$  AND  $\alpha=1.2, 1.5, 1.8$

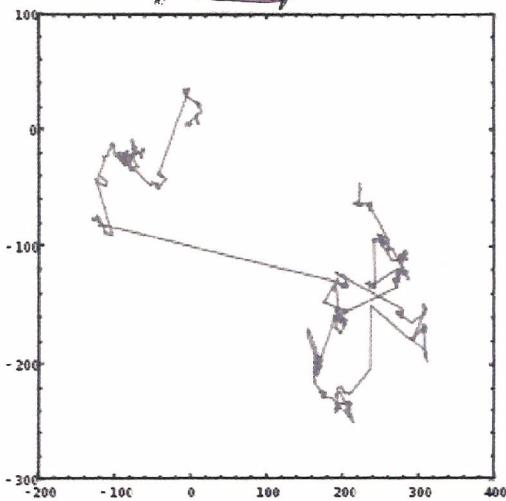
Levy

Figure 1. An example of 1000 steps of a Lévy flight in two dimensions. The origin of the motion is at [0,0], the angular direction is uniformly distributed and the step size is distributed according to a Lévy (i.e. stable) distribution with  $\alpha = 1$  and  $\beta = 0$  which is a Cauchy distribution. Note the presence of large jumps in location compared to the Brownian motion illustrated in Figure 2.

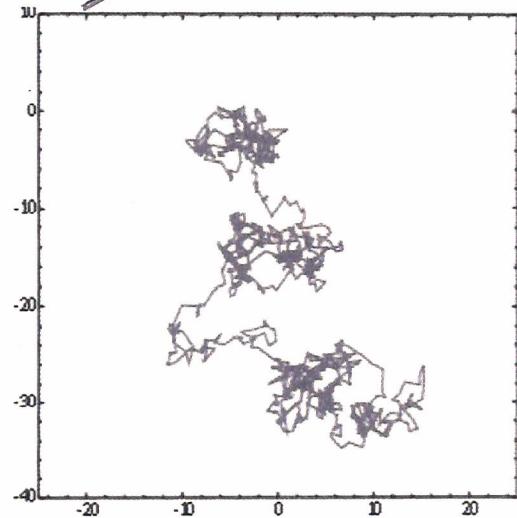
Brownian

Figure 2. An example of 1000 steps of an approximation to a Brownian motion type of Lévy flight in two dimensions. The origin of the motion is at [0, 0], the angular direction is uniformly distributed and the step size is distributed according to a Lévy (i.e. stable) distribution with  $\alpha = 2$  and  $\beta = 0$  (i.e., a normal distribution).

## Applications

The definition of a Lévy flight stems from the mathematics related to chaos theory and is useful in stochastic measurement and simulations for random or pseudo-random natural phenomena. Examples include earthquake data analysis, financial mathematics, cryptography, signals analysis as well as many applications in astronomy, biology, and physics.

Another application is the Lévy flight foraging hypothesis. When sharks and other ocean predators can't find food, they abandon Brownian motion, the random motion seen in swirling gas molecules, for Lévy flight — a mix of long trajectories and short, random movements found in turbulent fluids. Researchers analyzed over 12 million movements recorded over 5,700 days in 55 data-logger-tagged animals from 14 ocean predator species in the Atlantic and Pacific Oceans, including silky sharks, yellowfin tuna, blue marlin and swordfish. The data showed that Lévy flights interspersed with Brownian motion can describe the animals' hunting patterns.<sup>[7][8][9][10]</sup> Birds and other animals<sup>[11]</sup> (including humans)<sup>[12]</sup> follow paths that have been modeled using Lévy flight (e.g. when searching for food).<sup>[13]</sup> Biological flight data can also apparently be mimicked by other models such as composite correlated random walks, which grow across scales to converge on optimal Lévy walks.<sup>[14]</sup> Composite Brownian walks can be finely tuned to theoretically optimal Lévy walks but they are not as efficient as Lévy search across most landscapes types, suggesting selection pressure for Lévy walk characteristics is more likely than multi-scaled normal diffusive patterns.<sup>[15]</sup>

→ Class of L-stable processes  
is three "stable", as above,  
under addition.

includes:

Laplace → Gaussian

(only stable distribution  
with finite variance)

→ weak P-L laws with  $1 < \alpha < 2$   
(wild)

→ only possible limit laws of weighted sums of  
identical and ~~independent~~ random variables

→ density  $P$  of P-L laws (see p. 15)

$$G(b) = \int_{-\infty}^{\infty} e^{-bu} p(u) du$$

$$= \exp \left[ bu^* + M(b) \right]$$

and Laplace transform

$$\begin{cases} \alpha \\ u^* \\ M \rightarrow E(u) \end{cases}$$

→ Working principle:

- if - sum of many components  $\rightarrow$
- = Gaussian
- = skewed
- $E(u) < \infty$

→ reasonable assumption that  
follows P-L.